Actually, there will in general be several different operators of spin n+2 in the Wilson expansion, each making a contribution to the right hand side of Eq. (2), each with its own characteristic coefficient $b_n(i)$ and $a_n(i)$. For every n, it is the contribution with the smallest exponent that ultimately dominates at large q^2, and it is this contribution that is understood to be represented by the right hand side of Eq. (2).

Equation (2) describes the leading term in an expansion in inverse powers of log q^2 and q^2. It would be tempting to try to reconstruct the full structure function $F_i(\omega,q^2)$, for large q^2, by supposing that the correction terms in each moment are uniformly small, for all n, when q^2 exceeds some n-independent value. We shall in fact succumb to this temptation later on, but it is clear that any such procedure is highly speculative. At the present stage of theoretical understanding the only firm predictions that follow from the ideas of asymptotic freedom are those embodied in Eq. (2). Thus a sharp test of asymptotic freedom requires the difficult experimentation involved in extracting from the data the individual moments, as functions of q^2. A more modest experimental objective is the study of total neutrino and antineutrino cross sections as a function of energy. As is well known, strict scaling implies a linear growth with energy, at large energies, and indeed this kind of behavior is what is indicated by existing data. The question arises as to the growth properties that are to be expected for theories of the sort under present discussion. This is our first topic. We will see that both upper and lower bounds can be set on the growth rate, on the basis of the moment properties discussed above. It turns out that the bounds rather closely bracket a linear growth behavior.